

ON THE INVERSE RESONANCE PROBLEM

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I am reporting on joint work with

- Christer Bennewitz (Lund)
- Matthew Bledsoe (Birmingham, AL)
- Malcolm Brown (Cardiff)
- Ian Knowles (UAB)
- Marco Marletta (Cardiff)
- Sergey Naboko (St. Petersburg)
- Roman Shterenberg (UAB)

Introduction and History

“Über eine Frage der Eigenwerttheorie” (1928):

Wenn das Spektrum die Differentialgleichung wirklich vollständig definiert[e], so wäre es möglich, z. B. den Aufbau irgend eines Atomsystems praktisch aus dem Spektrum zu bestimmen, d. h. die Aufgabe zu lösen, welche sozusagen reziprok zum Schrödingerschen Problem steht.

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- Unless mentioned otherwise, we assume below a Dirichlet condition at zero.

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- Borg (1946) showed that, in general, two sets of eigenvalues are needed to identify a potential on an interval uniquely.
- Levinson (1949) and Marchenko (1950) used different sets of data: in addition to one set of (Dirichlet) eigenvalues one needs also either Neumann data of the eigenfunctions or the norming constants of the eigenfunctions.

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$$m(\lambda) = A\lambda + B + \int_{\mathbb{R}} \left(\frac{1}{t - \lambda} - \frac{t}{1 + t^2} \right) d\rho(t)$$

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- Gelfand-Levitan (1951): the spectral function ρ determines q uniquely.

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- Jost function: $\psi(\cdot, 0)$
- Marchenko (1955): eigenvalues, norming constants, and scattering phase ($2i\delta(k) = \overline{\psi(k, 0)}/\psi(k, 0)$) determine q uniquely.

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- If $\text{Im}(k) \leq 0$ and $\psi(k, 0) = 0$ then k^2 is a resonance. In this case $\psi(k, \cdot)$ satisfies the differential equation and the boundary condition at 0 but is not square integrable.
- Both are relevant/visible in spectroscopy.

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- It appears this went unnoticed for more than 40 years until Korotyaev (2000) and Zworski (2001/1988) pointed it out.

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- **Large resonances are physically insignificant.**
- Question: How can we state (and prove) this mathematically?

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Results

Stability for the inverse resonance theorem

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- Inside a disc of radius R all resonances and eigenvalues of q are ε -close to those of \tilde{q}

Then

$$\sup_{x \in [0, 1]} \left| \int_x^1 (q - \tilde{q}) dx \right| \leq f(\varepsilon, R)$$

where $f(\varepsilon, R) \rightarrow 0$ as $R \rightarrow \infty$ but $\varepsilon R^{1/6} \rightarrow 0$.

Stability in the case of a compact interval

- Ryabushko (1983): Suppose q_0 and q are real and have zero average. Then

$$\|q - q_0\|_{L^2} \leq C (\|\lambda(q) - \lambda(q_0)\|_{\ell^2} + \|\mu(q) - \mu(q_0)\|_{\ell^2})$$

where C depends on $\|q\|_2$ and $\|q_0\|_2$.

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- McLaughlin (1988) has a similar estimate involving one spectrum and norming constants.
- Marletta and myself (2005) gave an estimate (in terms of N and ε) on

$$\left| \int_0^x (q - q_0) dt \right| \leq f(\varepsilon, N)$$

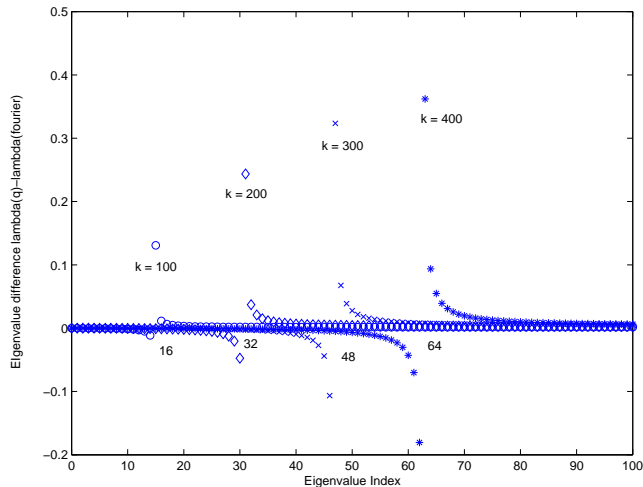
where $f(\varepsilon, N) \rightarrow 0$ as $N \rightarrow \infty$ but $\varepsilon \log N \rightarrow 0$ provided that $2N$ eigenvalues are known up to an error ε .

Comparison of eigenvalues

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- Marletta, Naboko, Shterenberg, W. (J. Anal. Math. 2011): Stability for several classes of Jacobi operators: Spectrum is (i) all of \mathbb{R} , (ii) a half-line, or (iii) one finite interval.

Full line problems

- Bledsoe (IEOT 2012): discrete case

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- Bledsoe (Inverse Problems 2012): continuous case
- Eigenvalues and resonances, i.e., the poles of the reflection coefficient, do not yet determine it.
- The zeros of the reflection coefficient are also needed.

Left-definite operators

- If $q \geq 0$ but no requirement on the sign of w is made one can develop a spectral and scattering theory for

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- The inverse resonance problem for the half-line case was treated by Bledsoe, W. (J. Math. Anal. Appl., 2015)

Outline of the proof

Transformation Operators

The Jost solutions associated with q and \tilde{q} are related by

$$\tilde{\psi}(z, x) = \psi(z, x) + \int_x^{2-x} K(x, t) \psi(z, t) dt$$

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$$\begin{aligned} K_{xx}(x, t) - K_{tt}(x, t) \\ = (\tilde{q}(x) - q(t))K(x, t) \end{aligned}$$

with the boundary conditions:

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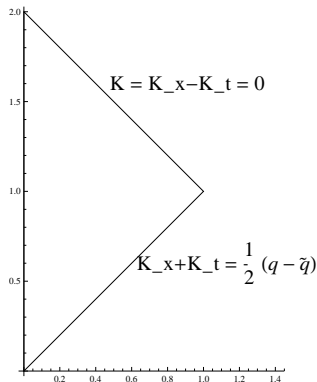
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with the boundary conditions:



Transformation Operators

The Jost solutions associated with q and \tilde{q} are related by

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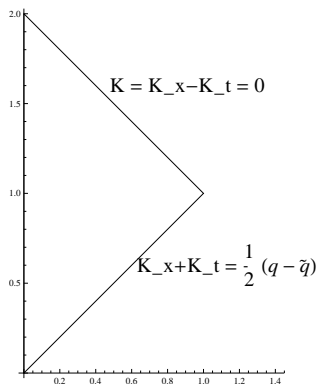
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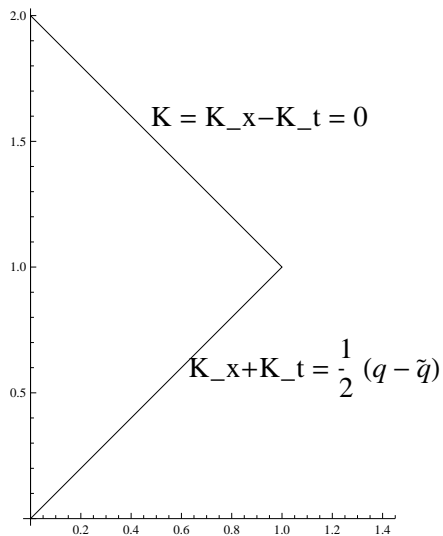
with the boundary conditions:

We need to estimate

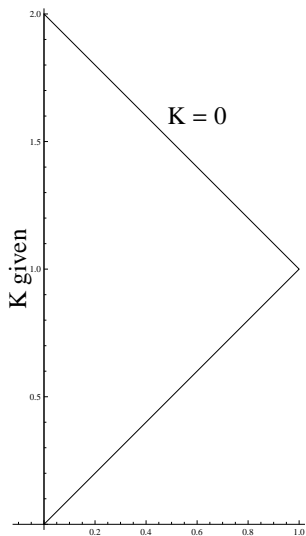
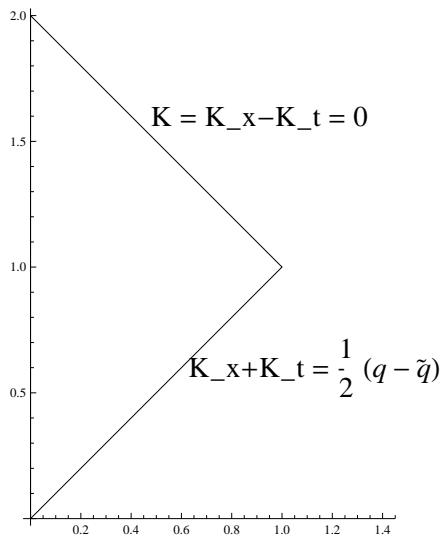
$$K(x, x) = \frac{1}{2} \int_x^1 (\tilde{q}(s) - q(s)) ds.$$



The wave equation



The wave equation



Solving the wave equation

The wave equation may be solved uniquely knowing $K(0, t)$, $0 \leq t \leq 2$ and the fact that $K(x, 2 - x) = 0$.

Iteration:

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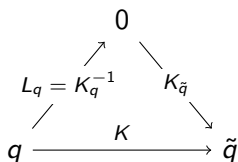
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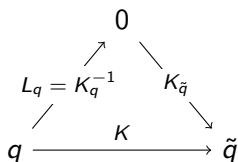
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Connecting with Jost functions I

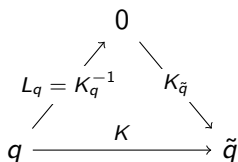


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- If f is entire of growth order at most one, then

$$f(z) = z^{n_0} e^{a+bz} \prod_{n=1}^{\infty} (1 - z/z_n) e^{z/z_n}.$$

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- $|\Pi(R, z) - 1| \leq C|z|^2/R$ when $2|z| < R$.
- This provides an estimate for $|z| < R^{1/3}$: $\psi(z, 0)/\tilde{\psi}(z, 0) \approx 1$ and hence

$$|\psi(z, 0) - \tilde{\psi}(z, 0)| \leq CR^{-1/3}.$$

Details

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Estimate for $|z| \geq R^{1/6}$

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- Here one needs the assumption that $\tilde{q} - q$ be in L^p .

Open problem

Two (or more) spectral bands

Suppose $q(x) = -2\wp(x + \omega)$. The spectrum of the associated Schrödinger operator has only absolutely continuous spectrum with one gap. All solutions of $-y'' + qy = \lambda y$ are explicitly known.

The inverse of the map $\wp(z) = \lambda$ maps the energy (λ) plane to parallelogram (the fundamental domain of \wp) in a one-to-two fashion. Compactly supported perturbations do not change the essential spectrum but introduce eigenvalues and resonances

Thank you for your attention!